Spectral geometry of tubes

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Why is the spectrum (of the Laplacian) so important ?

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We are thus lead to the study of the spectral-geometric problem for the Laplacian :

$$\begin{cases} -\Delta \psi = \lambda \psi & \text{in } \Omega, \\ \frac{\partial \psi}{\partial n} + \alpha \psi = 0 & \text{on } \partial \Omega. \end{cases}$$
 (Neumann $\alpha = 0$, Dirichlet $\alpha = \infty$)

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geometry of $\Omega \quad \leftrightarrows \quad \operatorname{spectrum}$ of $-\Delta$

But: spectrum known explicitly only for $\Omega = \mathbb{R}^d$, ball and parallelipiped \implies functional-analytic tools have to be employed

What is the spectrum?





 $\begin{array}{l} \textbf{i} \text{ How to correctly understand} \\ \begin{cases} -\Delta\psi=\lambda\psi & \text{in }\Omega, \\ \psi=0 & \text{on }\partial\Omega. \end{cases} \end{array} \begin{array}{l} \textbf{?} \end{array}$



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 \rightarrow Spectral problem for an unbounded (self-adjoint) operator in a Hilbert space:

$$-\Delta_D^{\Omega}: L^2(\Omega) \to L^2(\Omega): \left\{ \psi \mapsto -\Delta \psi \right\}$$
$$\mathsf{Dom}(-\Delta_D^{\Omega}):= \left\{ \psi \in W_0^{1,2}(\Omega) \mid \Delta \psi \in L^2(\Omega) \right\}$$



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NB (spectrum of an unbounded operator H) $\sigma(H) := \{ \lambda \in \mathbb{C} \mid H - \lambda I \text{ is not bijective} \}$



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NB (spectrum of an unbounded operator H) $\sigma(H) := \{ \lambda \in \mathbb{C} \mid H - \lambda I \text{ is not bijective} \}$ $= \underbrace{\sigma_{\rm disc}(H)}_{\bullet} \dot{\cup} \sigma_{\rm ess}(H)$

isolated eigenvalues of finite multiplicity

[Glazman 1963]

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• quasi-conical $:\iff \Omega \supset \{ \text{arbitrarily large balls} \}$



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• *quasi-bounded* : \iff neither q-conical nor q-cylindrical





Theorem.
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Proof. Weyl's criterion:

 $\lambda \in \sigma(H) \iff \exists \{\psi_n\} \subset \mathsf{Dom}(H) : \begin{cases} \|\psi_n\| = 1\\ \|(H - \lambda)\psi_n\| \to 0 \end{cases}$

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Given $k \in \mathbb{R}^d$ such that $|k|^2 = \lambda$, one takes

 $\psi_n(x) = e^{ik \cdot x} \tilde{\chi}_{B_n}(x) \text{ where: } \begin{cases} \{B_n\} = \text{a sequence of the enlarging balls} \\ \tilde{\chi}_B = \text{mollified and normalised } \chi_B \quad q.e.d. \end{cases}$

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Is there anything interesting to study ?

Theorem (Hardy inequality). Let $d \geq 3$. Then

$$\forall \psi \in W_0^{1,2}(\Omega), \quad \int_{\Omega} |\nabla \psi(x)|^2 \, \mathrm{d}x \geq \left(\frac{d-2}{2}\right)^2 \int_{\Omega} \frac{|\psi(x)|^2}{|x|^2} \, \mathrm{d}x$$



1877 - 1947

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1877–1947

Theorem (criticality of \mathbb{R}^1 and \mathbb{R}^2). Let d = 1, 2. For any non-positive measurable V,

 $\inf \sigma(-\Delta_D^{\mathbb{R}^d} + V) < 0$











not satisfied for *spiny urchin*

$$\sigma(-\Delta_D^{\Omega}) = \sigma_{\rm disc}(-\Delta_D^{\Omega})$$















Quasi-cylindrical domains



Quasi-cylindrical domains



distinguished subclass: TUBES

tubular neighbourhoods of submanifolds


Quasi-cylindrical domains



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ightarrow location of the essential spectrum

 \rightarrow existence of the discrete spectrum

ε





• ambient Riemannian manifold \mathbb{R}^d , $d \geq 2$

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• tube
$$\Omega := \left\{ x + \sum_{j=1}^{\operatorname{codim}\Sigma} t_j n_j(x) : (x,t) \in \Sigma \times \omega \right\}$$





 Σ_{ε}

Σ

 $\kappa < 0$

1n

 $\kappa = 0$

к>0

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Assumption. No self-intersections.

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- \rightarrow unbounded geometry
- \rightarrow uniform cross-section



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• compact submanifolds (with boundary),

but no systematic spectral-theoretic study of non-compact non-complete manifolds.

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GaAs/AlGaAs crescent shaped quantum wire

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The essential spectrum

Theorem ([D.K., Lu 2014 (J. Math. Phys.)]). If

• Σ is asymptotically flat (2nd fundamental form goes to 0 at infinity),

and

• the transport of ω along Σ is asymptotically parallel (relevant only if $\operatorname{codim} \Sigma \geq 2$),

then

$$\sigma_{\mathrm{ess}}(-\Delta_D^{\Omega}) = [E_1, \infty).$$

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$$\begin{array}{lll} \textit{Proof.} & -\Delta_D^{\Omega} & \simeq & -G^{\frac{1}{2}} \partial_i G^{\frac{1}{2}} G^{ij} \partial_j \\ & & L^2(\Omega) & & L^2(\Sigma \times \omega, \mathrm{d}x \, \mathrm{d}t) \end{array}$$



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and

Weyl's criterion adapted to quadratic forms:

$$\lambda \in \sigma(H) \iff \exists \{\psi_n\} \subset \underbrace{\mathsf{Dom}(H^{\frac{1}{2}})}_{\mathcal{H}_1} : \begin{cases} \|\psi_n\|_{\mathcal{H}} = 1\\ \|(H-\lambda)\psi_n\|_{\mathcal{H}_1^*} \to 0 \end{cases}$$
$$\mathsf{Dom}(H) \subset \mathcal{H}_1 \subset \mathcal{H} = \mathcal{H}^* \subset \mathcal{H}_1^*$$



Quantum layers: the discrete spectrum

 $\dim \Sigma = 2, \operatorname{codim} \Sigma = 1$

Theorem ([Duclos, Exner, D.K. 2001 (*Comm. Math. Phys.*)], [Carron, Exner, D.K. 2004 (*J. Math. Phys.*)]**).**

Let $K \in L^2(\Sigma)$ and $\Sigma \neq \mathbb{R}^2$. If

- $\int_{\Sigma} K \leq 0$, or
- ω is thin enough,

•
$$\int_{\Sigma} M^2 = \infty$$
 but $\nabla M \in L^2(\Sigma)$,

• $\Sigma \supset$ cylindrically symmetric end E with $\int_E K > 0$,

then

N

$$\inf \sigma(-\Delta_D^{\Omega}) < E_1.$$



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Let $K \in L^2(\Sigma)$ and $\Sigma \neq \mathbb{R}^2$. If Σ_{ε} In $\kappa < 0$ • $\int_{\Sigma} K \leq 0$, or Σ • ω is thin enough. $\kappa = 0$ or • $\int_{\Sigma} M^2 = \infty$ but $\nabla M \in L^2(\Sigma)$, $\kappa > 0$ or • $\Sigma \supset$ cylindrically symmetric end E with $\int_{F} K > 0$, then $\inf \sigma(-\Delta_D^{\Omega}) < E_1$.

Corollary. If Σ is asymptotically flat and any of the conditions above hold, then

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Proof. Test function $1 \times \mathcal{J}_1$ where \mathcal{J}_1 is the first eigenfunction of $-\Delta_D^{\omega}$, etc. q.e.d.

Quantum layers: examples





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Quantum layers: examples



 $\dim \Sigma = 2, \operatorname{codim} \Sigma = 1$

• extensions to higher dimensions and codimensions [Lin, Lu 2006], [Lin, Lu 2007]

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- extensions to tubes in curved ambient manifolds [D.K. 2003], [D.K. 2006], [Kolb, D.K. 2014], [Wachsmuth, Teufel 2013]



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- magnetic field $-\Delta \rightsquigarrow (-i\nabla A)^2$ [D.K., Raymond, Tušek 2015]



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Theorem ([D.K., Raymond, Tušek 2015 (*J. Geom. Anal.*)]**).** Replace $\omega \mapsto \varepsilon \omega$ with $\varepsilon > 0$. Then (dim $\Sigma = 2$, codim $\Sigma = 1$, A = 0)

$$-\Delta_D^{\Omega_{\varepsilon}} - \frac{E_1}{\varepsilon^2} \quad \xrightarrow{\text{n.r.s.}}_{\varepsilon \to 0} \quad -\Delta^{\Sigma} + K - M^2$$





 $\dim \Sigma = 1$, $\operatorname{codim} \Sigma = 2$

 $\Sigma := \{\Gamma(s) : s \in \mathbb{R}\}, \ \Gamma : \mathbb{R} \to \mathbb{R}^3 \text{ unit-speed } (|\dot{\Gamma}| = 1) \text{ immersion, curvature } \kappa := |\ddot{\Gamma}|$

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Frenet frame

$$\begin{pmatrix} \dot{\Gamma} \\ N \\ B \end{pmatrix} = \begin{pmatrix} 0 & \kappa & 0 \\ -\kappa & 0 & \tau \\ 0 & -\tau & 0 \end{pmatrix} \begin{pmatrix} \dot{\Gamma} \\ N \\ B \end{pmatrix}$$

 $\kappa > 0$

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Frenet frame versus $\begin{pmatrix} \dot{\Gamma} \\ N \\ B \end{pmatrix} = \begin{pmatrix} 0 & \kappa & 0 \\ -\kappa & 0 & \tau \\ 0 & -\tau & 0 \end{pmatrix} \begin{pmatrix} \dot{\Gamma} \\ N \\ B \end{pmatrix}$

 $\kappa > 0$

relatively parallel frame

$$\begin{pmatrix} \dot{\Gamma} \\ N_1 \\ N_2 \end{pmatrix}^{\bullet} = \begin{pmatrix} 0 & k_1 & k_2 \\ -k_1 & 0 & 0 \\ -k_2 & 0 & 0 \end{pmatrix} \begin{pmatrix} \dot{\Gamma} \\ N_1 \\ N_2 \end{pmatrix}$$
$$\kappa^2 = k_1^2 + k_2^2$$

[Bishop 1975] [D.K., Šediváková 2012]

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Quantum tubes: the geometry

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Theorem ([Chenaud, Duclos, Freitas, D.K. 2005 (Differential Geom. Appl.)]).

If $\kappa \neq 0$ and $\dot{\theta} = 0$, then

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bending acts as an *attractive* interacion



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Twisting versus bending

Corollary. Let $\dot{\theta} \neq 0$, $\dot{\theta} \in C_0(\mathbb{R})$ and ω is not circular. Then there exists $\epsilon > 0$,

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Theorem ([D.K., Šediváková 2012 (*Rev. Math. Phys.*)]**).** Replace $\omega \mapsto \varepsilon \omega$ with $\varepsilon > 0$. Then

$$-\Delta_D^{\Omega_{\varepsilon}} - \frac{E_1}{\varepsilon^2} \quad \xrightarrow[\varepsilon \to 0]{\text{n.r.s.}} \quad -\Delta^{\Sigma} - \frac{\kappa^2}{4} + \|\partial_{\tau}\mathcal{J}_1\|^2 \dot{\theta}^2$$

Remark. Previous related results: [Bouchitté, Mascarenhas, Trabucho 2007], [Wachsmuth, Teufel 2013], [de Oliveira 2010].

Application to the heat equation

$$\begin{cases} \frac{\partial u}{\partial t} - \Delta_x u = 0, \quad (x, t) \in \Omega \times (0, \infty), \\ u(x, 0) = u_0(x). \end{cases}$$



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twisting \implies faster cool down (death of a Brownian particle) in twisted tubes

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Open problem: ¿ Non-standard Weyl-type asymptotics in the case 2 ?